

Adaptivity and blowup detection for nonlinear evolution PDEs

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Based on joint work with:

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Overview

- Motivation – blowup detection
- Rigorous **a posteriori bounds** for semilinear parabolic problems, valid up to blowup time
- Adaptivity and **estimation of blowup time** and near blowup



A. CANGIANI, E. H. GEORGOULIS, I. KYZA, S. METCALFE, *Adaptivity and blow-up detection for nonlinear evolution problems*, in review.

- adaptive (high-order) methods for **non-polygonal interface** problems



A. CANGIANI, E. H. GEORGOULIS, Y. SABAWI, *Adaptive discontinuous Galerkin methods for non-polygonal interface problems*, in preparation.

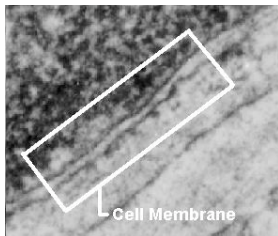
A “simple” test case

$u_t - \Delta u = u^2$, $u(0, x)$ Gaussian \rightsquigarrow single point blowup

Motivation

- blowup time estimation is interesting for applications (physical/chemical reactions, chemotaxis(?), etc.)
- a priori/analytical knowledge of blowup times is **known for few model problems only**
- a general enough error estimation framework could **provide insight** for interaction of non-linearities with other phenomena, such as advection, interfaces, etc.

Mass transfer through semi-permeable membranes



We consider a number of solutes subject to:

- **diffusion** and **advection** on both sides of the membrane
- **nonlinear reactions** with other solutes
- **mass transfer** across the membrane

(**red** = linear, **green** = nonlinear)

Cargo·Imp concentration in cell signal transduction

Simpler model problem...

Consider advection-diffusion-reaction PDE problem on a single domain $\Omega \subset \mathbb{R}^2$:

$$\begin{aligned}\partial_t u - \kappa \Delta u + \mathbf{a} \cdot \nabla u + f(u) &= 0 && \text{in } \Omega \times (0, T], \\ u &= 0 && \text{on } \partial\Omega \times (0, T], \\ u(\cdot, 0) &= u_0 && \text{in } \Omega,\end{aligned}$$

with $f(u) = -u^2$, $\kappa > 0$.

- analytical results on blowup times ?
- effect(s) of advection w.r.t. to blowup?
- inclusion of interfaces in the mix? (ongoing...)

For accessibility, in this talk, I shall mostly discuss the even simpler problem:

$$\begin{aligned}\partial_t u - \Delta u &= u^2 && \text{in } \Omega \times (0, T], \\ u &= 0 && \text{on } \partial\Omega \times (0, T], \\ u(\cdot, 0) &= u_0 && \text{in } \Omega,\end{aligned}$$

Aim

Estimation of **blowup time** & *space-time error control* near blowup

Blowup detection & error control

A case for **adaptivity**:

- **Extremely fine (space-time) local resolution needed** to approach blowup time
- “Standard” a priori error analysis & uniform meshes:
at time $t = T_{blowup} - \varepsilon$, constants $O(e^{1/\varepsilon})$ appear in bounds. $\varepsilon \rightarrow 0\dots$
- Problem of **missing the blowup!**

Dominant approach in the literature: rescaling/use of PDE ‘similarity’ properties to (r -)adapt/rescale discretisation parameters. Nakagawa ('75), Berger & Kohn ('89), Stuart & Floater ('90), Tourigny & Sanz-Serna ('92), Budd, Huang & Russell ('96) ...

Approach

Construction of adaptive algorithms via rigorous *a posteriori error estimates*

- **Limited literature** on a posteriori error control & adaptivity in this context.
Karakashian & Plexousakis ('96), Kyza & Makridakis ('11)
- **Conditional a posteriori error estimates:** final estimates **hold under some computationally verifiable conditions**

Step back to ODEs...

ODE initial value problem: find $u : [0, T] \rightarrow \mathbb{R}$ such that

$$\begin{aligned}\frac{du}{dt} &= f(u) := u^p, & \text{in } (0, T], \\ u(0) &= u_0,\end{aligned}$$

with $\mathbb{N} \ni p \geq 2$, so that the solution blows up in finite time, say T^* .

Three different one step schemes: set $U^0 := u_0$; for $k = 1, \dots, N$, solve for U^k :

$$\frac{U^k - U^{k-1}}{\tau_k} = F(U^{k-1}, U^k),$$

with F one of the following three classical approximations of f :

Explicit Euler $F(U^{k-1}, U^k) = f(U^{k-1}),$

Implicit Euler $F(U^{k-1}, U^k) = f(U^k),$

RK2/Improved Euler $F(U^{k-1}, U^k) = \frac{1}{2} (f(U^{k-1}) + f(U^{k-1} + \tau_k f(U^{k-1}))).$

An a posteriori error estimate

Let $U : [0, T] \rightarrow \mathbb{R}$ p/w linear interpolant of $\{U^k\}$ at t^k , viz

$$U(t) := \ell_{k-1}(t)U^{k-1} + \ell_k(t)U^k, \quad t \in (t^{k-1}, t^k],$$

Hence, on each interval $(t^{k-1}, t^k]$, we have

$$\frac{dU}{dt} = F(U^{k-1}, U^k).$$

Therefore, on each time interval $(t^{k-1}, t^k]$, the error $e := u - U$ satisfies

$$\frac{de}{dt} = f(u) - F(U^{k-1}, U^k) = f(U) + f'(u)e + \sum_{j=2}^p \frac{f^{(j)}(U)}{j!} e^j - F(U^{k-1}, U^k),$$

or, setting $\eta_k := f(U) - F(U^{k-1}, U^k)$, we have

$$\frac{de}{dt} = \eta_k + \left(f'(U) + \sum_{j=2}^p \frac{f^{(j)}(U)}{j!} e^{j-1} \right) e.$$

An a posteriori error estimate

Gronwall's inequality, therefore, implies

$$|e(t)| \leq H_k(t) G_k \phi_k,$$

$$\text{with } H_k(t) := \exp\left(\sum_{j=2}^p \int_{t^{k-1}}^t \frac{|f^{(j)}(U)|}{j!} |e|^{j-1} ds\right), \quad G_k := \exp\left(\int_{t^{k-1}}^{t^k} |f'(U)| ds\right),$$

and $\phi_k := |e(t^{k-1})| + \int_{t^{k-1}}^{t^k} |\eta_k| ds.$

Theorem (Conditional error estimate)

For $k = 1, \dots, N$, the following a posteriori estimate holds:

$$\max_{t \in [t^{k-1}, t^k]} |e(t)| \leq \delta_k G_k \phi_k,$$

provided that $\delta_k > 1$ is chosen so that

$$\sum_{j=2}^p (\delta_k G_k \phi_k)^{j-1} \int_{t^{k-1}}^{t^k} \frac{|f^{(j)}(U(s))|}{j!} ds - \log(\delta_k) = 0.$$

Two algorithms: 1

Algorithm 1 ODE Algorithm 1

1: **Input:** $f, F, u_0, \tau_1, \text{tol}$.
2: Compute U^1 from U^0 .
3: **while** $\int_{t^0}^{t^1} |\eta_1| ds > \text{tol}$ **do**
4: $\tau_1 \leftarrow \tau_1/2$.
5: Compute U^1 from U^0 .
6: **end while**
7: Compute δ_1 .
8: Set $k = 0$.
9: **while** δ_{k+1} exists **do**
10: $k \leftarrow k + 1$.
11: $\tau_{k+1} = \tau_k$.
12: Compute U^{k+1} from U^k .
13: **while** $\int_{t^k}^{t^{k+1}} |\eta_{k+1}| ds > \text{tol}$ **do**
14: $\tau_{k+1} \leftarrow \tau_{k+1}/2$.
15: Compute U^{k+1} from U^k .
16: **end while**
17: Compute δ_{k+1} .
18: **end while**
19: **Output:** k, t^k .

$$\eta_k := f(U) - F(U^{k-1}, U^k)$$

Absolute tolerance: tol

$$|T_{\text{blowup}} - T_{\text{final}}| \sim N^{-r}$$

Method	$p = 2$	$p = 3$
Implicit Euler	$r \approx 0.66$	$r \approx 0.79$
Explicit Euler	$r \approx 1.35$	$r \approx 1.60$
Improved Euler	$r \approx 1.2$	$r \approx 1.48$

Two algorithms: 2

Algorithm 2 ODE Algorithm 2

```
1: Input:  $f, F, u_0, \tau_1, \text{tol}$ .
2: Compute  $U^1$  from  $U^0$ .
3: while  $\int_{t^0}^{t^1} |\eta_1| ds > \text{tol}$  do
4:    $\tau_1 \leftarrow \tau_1/2$ .
5:   Compute  $U^1$  from  $U^0$ .
6: end while
7: Compute  $\delta_1$ .
8:  $\text{tol} = G_1 * \text{tol}$ .
9: Set  $k = 0$ .
10: while  $\delta_{k+1}$  exists do
11:    $k \leftarrow k + 1$ .
12:    $\tau_{k+1} = \tau_k$ .
13:   Compute  $U^{k+1}$  from  $U^k$ .
14:   while  $\int_{t^k}^{t^{k+1}} |\eta_{k+1}| ds > \text{tol}$  do
15:      $\tau_{k+1} \leftarrow \tau_{k+1}/2$ .
16:     Compute  $U^{k+1}$  from  $U^k$ .
17:   end while
18:   Compute  $\delta_{k+1}$ .
19:    $\text{tol} = G_{k+1} * \text{tol}$ .
20: end while
21: Output:  $k, t^k$ .
```

$$\eta_k := f(U) - F(U^{k-1}, U^k)$$

Relative tolerance: $G_{k+1} * \text{tol}$

$$|T_{\text{blowup}} - T_{\text{final}}| \sim N^{-r}$$

Method	$p = 2$	$p = 3$
Implicit Euler	$r \approx 1.00$	$r \approx 1.00$
Explicit Euler	$r \approx 1.45$	$r \approx 1.43$
Improved Euler	$r \approx 2.03$	$r \approx 2.03$

Back to PDEs: time semi-discretisation

$0 := t_0 < t_1 < \dots < t_N := T$ partition of $[0, T]$, $\tau_k := t_{k+1} - t_k$,

Implicit-Explicit (IMEX) Euler method: find $U^k \in H_0^1(\Omega)$, $k = 0, 1, \dots, N - 1$:

$$\frac{U^{k+1} - U^k}{\tau_k} - \Delta U^{k+1} = f(U^k), \text{ with } U^0 = u_0$$

Why IMEX

- implicit on diffusion \Rightarrow stability
- explicit on nonlinear reaction \Rightarrow advantageous approximation near blowup

Error equation

Let $U : [0, T] \rightarrow H_0^1(\Omega)$ linear interpolant of $\{U^k\}_k$.

Let $e := u - U$. Then, for $f(u) = u^2$, we have

$$\partial_t e - \Delta e = 2Ue + e^2 + r_{k+1}$$

in $(t_k, t_{k+1}]$ with $r_{k+1} := [f(U) - f(U^k)] + (t_{k+1} - t)(U^{k+1} - U^k)/\tau_k$.

Energy estimate:

$$\frac{d}{dt} \|e(t)\|^2 + \|\nabla e(t)\|^2 \leq 4\|U(t)\|_{L^\infty} \|e(t)\|^2 + 2\langle e^2, e \rangle + \|r_{k+1}(t)\|_{-1}^2$$

Using $\langle e^2, e \rangle \leq \|e\|_{L^\infty} \|e\|^2$, Gronwall's inequality gives

$$\max_{0 \leq t \leq T} \|e(t)\|^2 \leq \exp\left(2 \int_0^T [2\|U(t)\|_{L^\infty} + \|e(t)\|_{L^\infty}] dt\right) \sum_{k=0}^{N-1} \int_{t_k}^{t_{k+1}} \|r_{k+1}(t)\|_{-1}^2 dt$$

- Is a fully a posteriori bound possible from this?
- Behaviour of the constant in the run up to blowup?

Exponent growth (Kyza & Makridakis ('11))

Gronwall exponent:

$$\int_0^T 2\|U(t)\|_{L^\infty} + \|e(t)\|_{L^\infty} dt$$

On run up to blowup time: $\|u(t)\|_{L^\infty} \sim \frac{1}{(T_{blowup} - t)}$ (Merle & Zaag ('00))

We can infer that we also have, approximately,

$$\|U(t)\|_{L^\infty}, \|e(t)\|_{L^\infty} \sim \frac{1}{(T_{blowup} - t)}$$

Hence, at $T := T_{blowup} - \varepsilon$, $\varepsilon > 0$, the exponent scales like

$$\int_0^{T_{blowup} - \varepsilon} \frac{1}{T_{blowup} - t} dt = \ln\left(\frac{T_{blowup}}{\varepsilon}\right)$$

and, thus, the Gronwall constant scales like

$$\exp\left(\int_0^{T_{blowup} - \varepsilon} \frac{1}{T_{blowup} - t} dt\right) \sim \frac{C(T_{blowup})}{\varepsilon^q}$$

i.e., **polynomial growth w.r.t ε !** Challenge: Leads to practical algorithm?

Time-discrete scheme a posteriori bound of (Kyza & Makridakis ('11))

fixed point arguments + semigroup theory \implies a posteriori estimate

Conditional a posteriori error estimates

- $L^\infty(L^\infty)$: $\|e\|_{L^\infty(L^\infty)} \leq e^{1/8+4 \int_0^T \|U(s)\|_{L^\infty} ds} \sum_{k=0}^{N-1} \int_{t_k}^{t_{k+1}} \|r_{k+1}(s)\|_{L^\infty} ds$
- $L^\infty(L^2)$: $\max_{0 \leq t \leq T} \|e(t)\|^2 \leq e^{1/8+4 \int_0^T \|U(s)\|_{L^\infty} ds} \sum_{k=0}^{N-1} \int_{t_k}^{t_{k+1}} \|r_{k+1}(s)\|_{-1}^2 ds$

provided τ_k are chosen so that

$$e^{4 \int_0^T \|U(s)\|_{L^\infty} ds} \sum_{k=0}^{N-1} \int_{t_k}^{t_{k+1}} \|r_{k+1}(s)\|_{L^\infty} ds \leq \frac{3}{16} \rho, \quad \text{with } \rho < \frac{1}{16T}$$

Global condition! efficient time-adaptive algorithm? – More local conditions?

- For fully-discrete L^∞ -norm a posteriori bounds for resp. elliptic problem are needed in this framework.
- growth range u^p , $p > 1$.

A new a posteriori bound (fully discrete bound also available)

New conditional estimate

$$\max_{t_k \leq t \leq t_{k+1}} \|e(t)\|^2 \leq \delta_{k+1} e^{4 \int_{t_k}^{t_{k+1}} \|U(s)\|_{L^\infty} ds} \left[\|e(t_k)\|^2 + \int_{t_k}^{t_{k+1}} \|r_{k+1}(s)\|_{-1}^2 ds \right]$$

with $e(0) = 0$, where $\delta_{k+1} > 1$ is chosen to satisfy

$$\beta^2 \delta_{k+1} e^{4 \int_{t_k}^{t_{k+1}} \|U(s)\|_{L^\infty} ds} \left[\|e(t_k)\|^2 + \int_{t_k}^{t_{k+1}} \|r_{k+1}(s)\|_{-1}^2 ds \right] \tau_k - \ln \delta_{k+1} = 0$$

Ingredients:

- Gagliardo–Nirenberg inequality: $\langle e^2, e \rangle(t) \leq \beta \|e(t)\|^2 \|\nabla e(t)\|$
- a new **local on each time-step** continuation argument
(“global” continuation arguments in this spirit: [Kessler, Nochetto & Schmidt \('04\)](#), [Bartels \('05\)](#), [Cangiani, G. & Jensen \('13\)](#))

Condition – discussion

Condition satisfied only if

$$\beta^2 e^{4 \int_{t_k}^{t_{k+1}} \|U(s)\|_{L^\infty} ds} \left[\|e(t_k)\|^2 + \int_{t_k}^{t_{k+1}} \|r_{k+1}(s)\|_{-1}^2 ds \right] \tau_k < \frac{1}{e}$$

This implies *restriction on time-steps* τ_k , i.e., *conditional estimates*.

Is condition practical?

- Say $t_{k+1} = T_{blowup} - \varepsilon$, then $e^{4 \int_{t_k}^{t_{k+1}} \|U(s)\|_{L^\infty} ds} \sim \left(\frac{\tau_k + \varepsilon}{\varepsilon} \right)^q$
- $\int_{t_n}^{t_{n+1}} \|r_{n+1}(s)\|_{-1}^2 ds$ estimable via a posteriori bounds

δ_{n+1} computed via, e.g., Newton's method.

Semilinear advection-diffusion problem with blowup

$$\begin{aligned}u_t - \epsilon \Delta u + \mathbf{a} \cdot \nabla u &= u^2 + g && \text{in } \Omega \times (0, T], \\u &= 0 && \text{on } \partial\Omega \times (0, T], \\u(\cdot, 0) &= u_0 && \text{in } \Omega,\end{aligned}$$

Spatial discretisation: **Discontinuous Galerkin (dG) method** with upwind flux.

dG IMEX method: find $U_h^{k+1} \in \mathbb{V}_h^{k+1}$ s.t.

$$\left\langle \frac{U_h^{k+1} - U_h^k}{\tau_k}, V_h \right\rangle + B(t_{k+1}; U_h^{k+1}, V_h) = \langle f(U_h^k) + g^k, V_h \rangle \quad \forall V_h \in \mathbb{V}_h^{k+1}$$

Continuation Argument + Elliptic Reconstruction



Conditional a posteriori error estimate in the $L^\infty(L^2)$ -norm for dG IMEX

For elliptic reconstruction see [Makridakis & Nochetto \('03\)](#), [Lakkis & Makridakis \('06\)](#), [G. Lakkis & Virtanen \('11\)](#), [Cangiani, G., & Metcalfe \('14\)](#)

Key attribute of the approach: **Flexibility on the elliptic operator!**

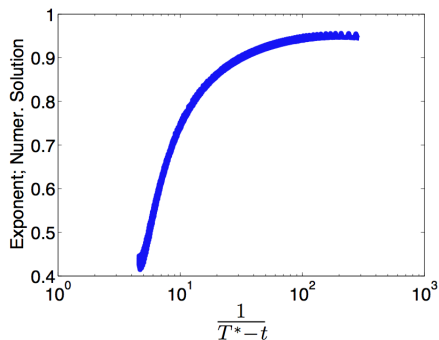
A simple test case: reaction-diffusion problem

$$\Omega = (-4, 4) \times (-4, 4), u_0(x, y) = 10e^{-2(x^2+y^2)} \rightsquigarrow \text{single point blowup}$$

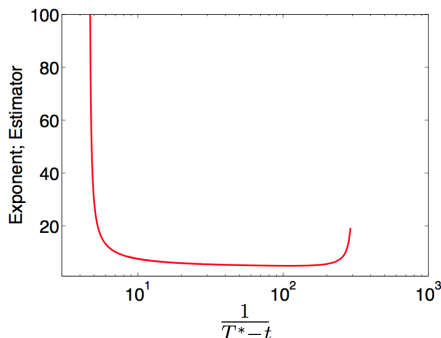
- space-time adaptive algorithm
- when condition fails, restarts with smaller timestep

A simple test case

tol	# time-steps	T	$\ U_h(T)\ _{L^\infty}$	Estimator
$(0.125)^{10}$	6956	0.21228	238.705	33426.7
$(0.125)^{11}$	14008	0.21375	343.078	36375.0
$(0.125)^{12}$	28151	0.21478	496.885	66012.8
$(0.125)^{13}$	35580	0.21549	722.884	157300.0



$$\|U_h(t)\|_{L^\infty} \sim \frac{1}{(T^* - t)^{p_k}}$$



plots depict: p_k vs. $\frac{1}{T^* - t}$

Observed rates of convergence

- $\|U_h(t)\|_{L^\infty}$ appears to blow up at the expected rate on the run-up to T^*
- we have $|T^* - T| \sim N^{-1/2}$
 - ▶ shortcoming of energy method?
 - ▶ semigroup techniques?
- Estimator blows up at faster rate, but delivers optimal blowup rate for the numerical solution regardless!

Numerical experiment – advection-diffusion problem

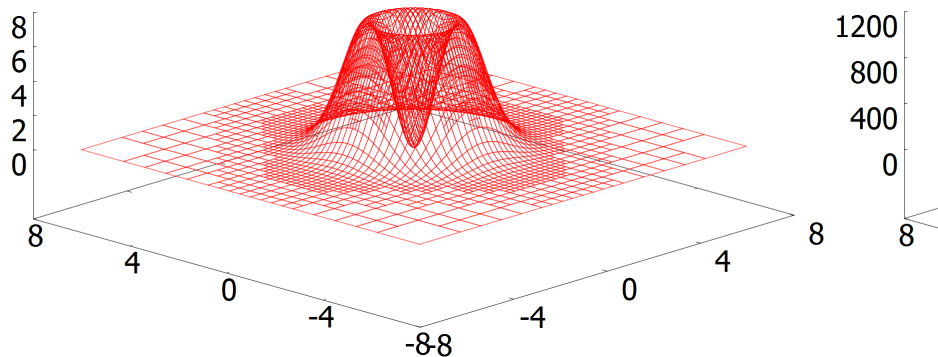
$$\Omega = (-4, 4) \times (-4, 4), \kappa = 1, \mathbf{a} \equiv (1, 1)^T, g \equiv 1, u_0 \equiv 0$$

tol	# time-steps	T	$\ U_h(T)\ _{L^\infty}$
1	4	0.78125	0.886245
0.125	10	0.976562	1.32178
$(0.125)^2$	54	1.31836	3.26904
$(0.125)^3$	119	1.41602	5.10672
$(0.125)^4$	252	1.48163	8.05863
$(0.125)^5$	520	1.51711	11.8193
$(0.125)^6$	1064	1.54467	18.1385
$(0.125)^7$	2158	1.56224	27.4045
$(0.125)^8$	4354	1.57402	41.3737
$(0.125)^9$	8792	1.58243	64.4503
$(0.125)^{10}$	17713	1.58770	99.1902
$(0.125)^{11}$	35580	1.59092	145.785
$(0.125)^{12}$	71352	1.59299	211.278

Numerical experiment – blowup on 1D manifold

$$\Omega = (-8, 8)^2, \kappa = 1, \mathbf{a} = (0, 0)^T, f_0 = 0$$

'volcano' type initial condition be given by $u_0 = 10(x^2 + y^2) \exp(-(x^2 + y^2)/2)$



Numerical experiment – blowup on 1D manifold

$$\Omega = (-8, 8)^2, \kappa = 1, \mathbf{a} = (0, 0)^T, f_0 = 0$$

'volcano' type initial condition be given by $u_0 = 10(x^2 + y^2) \exp(-(x^2 + y^2)/2)$

ttol^+	Time Steps	Estimator	Final Time	$\ U_h(T)\ _{L^\infty(\Omega)}$
8	3	15	0.06250	10.371
1	10	63	0.09375	14.194
0.125	36	211	0.11979	21.842
$(0.125)^2$	86	533	0.13412	31.446
$(0.125)^3$	190	971	0.14388	45.122
$(0.125)^4$	404	1358	0.15072	64.907
$(0.125)^5$	880	5853	0.15601	98.048
$(0.125)^6$	1853	10654	0.15942	146.162
$(0.125)^7$	3831	21301	0.16176	219.423
$(0.125)^8$	7851	143989	0.16336	332.849
$(0.125)^9$	16137	287420	0.16442	505.236
$(0.125)^{10}$	32846	331848	0.16512	769.652
$(0.125)^{11}$	66442	626522	0.16558	1175.21

Notation

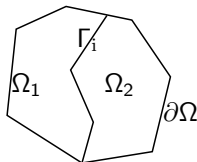
Consider an open polygonal domain $\Omega \subset \mathbb{R}^d$ subdivided into two subdomains Ω_1 and Ω_2 :

$$\Omega = \Omega_1 \cup \Omega_2 \cup \Gamma_i$$

$$\Gamma_i = \bar{\Omega}_1 \cap \bar{\Omega}_2$$

and set $\mathcal{H}^1 := [H^1(\Omega_1 \cup \Omega_2)]^n$, $n \in \mathbb{N}$.

PDE system:



$$\mathbf{u}_t - \nabla \cdot (A\nabla\mathbf{u} - U\mathbf{B}) + \mathbf{F}(\mathbf{u}) = \mathbf{0} \quad \text{in } (0, T] \times (\Omega_1 \cup \Omega_2),$$

$$\mathbf{u}(0, x) = \mathbf{u}_0(x) \quad \text{on } \{0\} \times \Omega,$$

$$\mathbf{u} = \mathbf{g}_D \quad \text{on } \Gamma_D,$$

$$(A\nabla\mathbf{u} - \chi^- U\mathbf{B})\mathbf{n} = \mathbf{g}_N \quad \text{on } \Gamma_N,$$

where χ^- the (vector-valued) characteristic function of the inflow part of $\partial\Omega$ and $U := \text{diag}(\mathbf{u})$. On Γ_i we impose:

$$(A\nabla\mathbf{u} - U\mathbf{B})\mathbf{n}|_{\Omega^1} = \mathbf{P}(\mathbf{u})(\mathbf{u}^2 - \mathbf{u}^1) - \{U\}_w R\mathbf{B}\mathbf{n}^1$$

$$(A\nabla\mathbf{u} - U\mathbf{B})\mathbf{n}|_{\Omega^2} = \mathbf{P}(\mathbf{u})(\mathbf{u}^1 - \mathbf{u}^2) - \{U\}_w R\mathbf{B}\mathbf{n}^2$$

Reflective membranes – time-dependent sharp features

uniform

adaptive

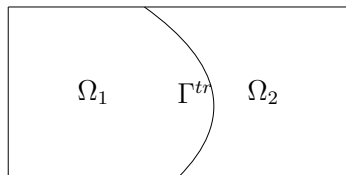
- a DG method & a priori analysis with solution boundedness assumption

Cangiani, G., Jensen ('13)

- No analytical information – numerics?

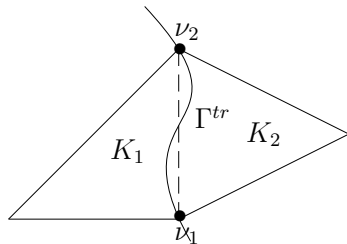
A posteriori error bounds for dG on curved geometries

$\Omega \subset \mathbb{R}^d$, $d = 2, 3$. $\Omega = \Omega_1 \cup \Omega_2 \cup \Gamma^{tr}$, with
 $\Gamma^{tr} := (\partial\Omega_1 \cap \partial\Omega_2) \setminus \partial\Omega$ Lipschitz.



$$\begin{aligned} -\Delta u &= f, & \text{in } \Omega_1 \cup \Omega_2, \\ u &= 0, & \text{on } \partial\Omega, \\ \mathbf{n}^1 \cdot \nabla u_1 &= C_{tr}(u_2 - u_1)|_{\Omega_1} & \text{on } \bar{\Omega}_1 \cap \Gamma^{tr}, \\ \mathbf{n}^2 \cdot \nabla u_2 &= C_{tr}(u_1 - u_2)|_{\Omega_2} & \text{on } \bar{\Omega}_2 \cap \Gamma^{tr}, \end{aligned}$$







where $u_i = u|_{\bar{\Omega}_i \cap \Gamma^{tr}}$, $i = 1, 2$, C_{tr} a given permeability constant.



Conclusions

- 'Energy methods' for error control are relevant and competitive for nonlinear evolution problems
- Error control for curved interfaces/boundaries

Some references

-  A. CANGIANI, E. H. GEORGOULIS, M. JENSEN, *Discontinuous Galerkin methods mass-transfer through semi-permeable membranes*, SIAM J. Numer. Anal., 51, pp. 2911-2934 (2013).
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Computation of numerical blowup time and blowup rate

Let t^* denote a **numerical blowup time**. We implement a set numerical experiments (**corresponding to different tolerances**) producing U_h^ℓ , $\ell = 1, \dots, L$ approximations to the exact solution u . Assume that

$$\|U_h^\ell(t_n)\|_{L^\infty} \sim \left(\frac{1}{t^* - t_n}\right)^p$$

- 1 Since $f(u) = u^2$, assume that $p = 1$, and use $\|U_h^L(t_{N-1})\|_{L^\infty}$, $\|U_h^L(T)\|_{L^\infty}$ to calculate

$$\left. \begin{aligned} \|U_h^L(t_{N-1})\|_{L^\infty} &= C_L \frac{1}{t^* - t_{N-1}} \\ \|U_h^L(T)\|_{L^\infty} &= C_L \frac{1}{t^* - T} \end{aligned} \right\} \Rightarrow t^* = \frac{T \|U_h^L(T)\|_{L^\infty} - t_{N-1} \|U_h^L(t_{N-1})\|_{L^\infty}}{\|U_h^L(t_{N-1})\|_{L^\infty} - \|U_h^L(T)\|_{L^\infty}}$$

- For the considered example, $t^* = 0.21705$
- 2 Consider $t^* (= 0.21705)$ as the **numerical blowup time**. We use $\|U_h^\ell(t)\|_{L^\infty}$, $\ell \neq L$ to compute the numerical blowup time:

$$p_n := \frac{\ln \left(\|U_h^\ell(t_{n+1})\|_{L^\infty} / \|U_h^\ell(t_n)\|_{L^\infty} \right)}{\ln \left((t^* - t_n) / (t^* - t_{n+1}) \right)}$$

- We expect $p_n \rightarrow 1$ as $n \rightarrow N$, for the considered model problem